

Let's think about entrainment:

$$\frac{DMS\bar{E}_c}{Dt} = Q_e - \frac{1}{m} \left(\frac{Dm}{Dt} \right)_e (MS\bar{E}_c - MS\bar{E}_e)$$

diabatic sources/sinks of MS \bar{E}

Dilution b.c. MS $\bar{E}_e < MS\bar{E}_c$
s.c. cloud is saturated

Mass exchange rate due to entrainment

Let's assume that $Q_e = 0$ and that

$$\frac{DMS\bar{E}_c}{Dt} = \frac{\partial MS\bar{E}_c}{\partial t} + \vec{v} \cdot \nabla MS\bar{E}_c + w \frac{\partial MS\bar{E}_c}{\partial z}$$

MS \bar{E} varies more rapidly in the vertical, even when accounting for $w \ll |\vec{v}|$

Expanding into MS \bar{E} budget.

$$\frac{\partial MS\bar{E}_c}{\partial z} = - \frac{1}{m} \left(\frac{\partial m}{\partial z} \right)_e (MS\bar{E}_c - MS\bar{E}_e)$$

$$e = \frac{1}{m} \left(\frac{\partial m}{\partial z} \right)_e = \text{entrainment rate}$$

$$\frac{\partial MS\bar{E}_c}{\partial z} = -e (MS\bar{E}_c - MS\bar{E}_e)$$

Entraining plume MS \bar{E} budget.

What if $T_e = T_c$ Convective Q \bar{E}
The cloud is saturated so $MS\bar{E}_c = MS\bar{E}_c^*$

$$\frac{\partial MS\bar{E}_c}{\partial z} = -e L_w (q_c^* - q_e)$$

↑ Dilution is related to the saturation deficit. $q^{\dagger} = q^* - q$

$$\beta = -\frac{\rho'}{\rho} g \approx \frac{T_c - T_e}{T} = \frac{T'}{T} = \frac{dT}{T}$$

$$dT \propto dMSE^*$$

$$\begin{aligned} dMSE^* &= d\phi T + d\Phi + L v dq^* \\ &= \phi dT + L v \frac{dq^*}{dT} dT \end{aligned}$$

$$dMSE^* = \phi \left(1 + \frac{L v \frac{dq^*}{dT}}{\phi} \right) dT$$

K_r

$$dT = \frac{1}{\phi K_r} dMSE^*$$

Plug this differential onto β :

$$\beta = \frac{MSE_c^* - MSE_e^*}{K_r \phi T}$$

Let's integrate the plume MSE budget

$$\int_{z_B}^z \frac{\partial MSE_c^*}{\partial z} dz = \int_{z_B}^z -\epsilon L v (q^* - q_e) dz$$

$$MSE_c^*(z) - MSE_c^*(z_B) = - \int_{z_B}^z \epsilon L v (q^* - q) dz$$

$$MSE_c^*(z) - MSE_B = - \int_{z_B}^z \epsilon L v (q^* - q) dz$$

Plugging this to β :

$$\beta = g \frac{MSE_B - MSE^*}{\phi K_r T} - \frac{g}{\phi K_r T} \int_{z_B}^z \epsilon L v (q^* - q) dz$$